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ESSnuSB near and far detector technology / strategy



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ESSnuSB workshop, Hamburg, 8 Oct 2020

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CP violation in ESSnuSB

 $P_{\mu \to e} \neq P_{\overline{\mu} \to \overline{e}}$

We will study ν_e and $\overline{\nu}_e$ appearance in ν_μ and $\overline{\nu}_\mu$ beam, respectively

The plan:

- 1. Run with v_{μ} and look at v_{e} appearance, then
- 2. Run with \overline{v}_{μ} and look at \overline{v}_{e} appearance

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ESSvSB v energy distribution (without optimisation)



- almost pure v_{μ} beam
- small v_e contamination which could be used to measure v_e crosssections in a near detector

	positive		negative	
	$N_{ u}~(imes 10^{10})/{ m m}^2$	%	$N_{ u}~(imes 10^{10})/{ m m}^2$	%
$ u_{\mu}$	396	97.9	11	1.6
$\bar{ u}_{\mu}$	6.6	1.6	206	94.5
ν_e	1.9	0.5	0.04	0.01
$\bar{\nu}_e$	0.02	0.005	1.1	0.5

at 100 km from the target and per year (in absence of oscillations)

(Nucl. Phys. B 885 (2014) 127)

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Purpose of the ESSnuSB detectors

Far detectors

•Observe ${}^{'}\!\bar{\nu}_{e}^{'}$ appearance in the ${}^{'}\!\bar{\nu}_{\mu}^{'} \to {}^{'}\!\bar{\nu}_{e}^{'}$ oscillation channel

Near detectors

- Constrain the prompt neutrino flux
- Measure neutrino interaction cross-sections (both inclusive and exclusive)



Neutrino baseline



• Garpenberg mine, 540 km from the neutrino source, corresponding to 2nd oscillation maximum.

Alternatives:

- Zinkgruvan mine, 340 km from source
- Garpenberg and Zinkgruvan, 250 kt each





Far detectors

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Far detectors



Baseline – two identical modules.

• It would be cheaper to have one larger module, but then we would have a problem with statics of the cave

Each module is a standing cylinder:

- diameter D = 78 m, height h = 78 m
 - D = h minimizes surface/volume ratio
- 373k m³ total volume
- 270k m³ fiducial volume (~10xSuperK)
- Readout: 38k 20" PMTs
- 30% optical coverage

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Possible positions at Zinkgruvan mine



Expected interaction rates in Far Detectors

Plots and numbers by L. Halić

Neutrino mode

Antineutrino mode



Approx. expected number of interactions at 540 km in 540 kt of water for 2.16 x 10²³ p.o.t. (effective year), assuming $\delta_{CP} = 0$:

Channel	Expected number
$\nu_{\mu} \rightarrow \nu_{e}$	200
$\nu_{\mu} \rightarrow \nu_{\mu}$	3600
$\nu_e \rightarrow \nu_e$	30

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Channel	Expected number
$\overline{\nu}_{\!\mu} \to \! \overline{\nu}_{\!e}$	40
$\overline{\nu}_{\!\mu} \to \overline{\nu}_{\!\mu}$	600
$\overline{\nu_{\!_e}} \rightarrow \overline{\nu_{\!_e}}$	3





Neutrino energy reconstruction

Kinematical neutrino energy reconstruction formula

$$E_{\nu}^{rec} = \frac{m_f^2 - (m_i')^2 - m_l^2 + 2m_i'E_l}{2(m_i' - E_l + p_l\cos\theta_l)}$$
(4)

where E_{ν}^{rec} is the reconstructed neutrino energy, m_i and m_f are the initial and final nucleon masses respectively, and $m'_i = m_i - E_{\rm b}$, where $E_{\rm b} = 27 \,\text{MeV}$ is the binding energy of a nucleon inside ¹⁶O nuclei. E_l , p_l and θ_l are the reconstructed lepton energy, momentum, and angle with respect to the beam, respectively. The selec-

From: Phys. Rev. D 96, 092006

Given that you know:

- magnitude of the outgoing charged lepton momentum
- its angle w.r.t. incoming neutrino
- that it is a quasielastic QES interaction
- which nucleus neutrino interacted with (e.g. oxygen-16)
 you can **approximately** calculate neutrino energy.

Intrinsic uncertainties come from nuclear effects, most notably **fermi motion** of nucleons in nuclei.







Conclusion: with a perfect detector using kinematical reconstruction formula we get energy resolution of about **110 MeV for neutrinos** and **65 MeV for antineutrinos**.

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S Difference between reconstructed and true lepton momentum





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Difference between reconstructed and true v energy





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CSI





Far detector conclusions / strategy

- Main purpose: observe $\stackrel{(i)}{\nu}_{e}$ appearance in the $\stackrel{(i)}{\nu}_{\mu} \rightarrow \stackrel{(i)}{\nu}_{e}$ oscillation channel
- Energy reconstruction:
 - Using the kinematical reconstruction formula, there is an intristic resolution of **110 MeV for neutrinos** and **65 MeV for antineutrinos**.
 - Therefore, trying to improve the resolution of charged lepton momentum produces very **diminishing returns** – no reason to improve it further
 - Having calorimetry could significantly improve energy resolution
 - but, you can't have calorimetry in pure WC detector
- Things that we are optimizing (but carefully in order not to screw up nonbeam physics potential):
 - event tagging (v_{μ} CC, v_{e} CC, NC discrimination) work in progress
 - detector geometry w.r.t. the local rock conditions
 - PMT coverage





Near detectors

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Near Detectors



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Neutrino mode

Antineutrino mode



Approximate expected interactions at 250 m in 500 t of water for 2.16 x 10²³ p.o.t. (effective year):

Neutrino	Expected number
ν_{μ}	27.5 M
$\overline{\nu}_{\mu}$	66 k
ν _e	150 k
$\overline{v_{e}}$	300

Neutrino	Expected number
ν_{μ}	265 k
$\overline{\nu}_{\mu}$	4.7 M
v _e	1.8 k
$\overline{\nu}_{e}$	15 k

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ND design discussion







Near SFGD detector



• Baseline dimensions:

- 140 x 140 x 50 cubes
- each cube 1cm x 1cm x 1cm
- Possible to do calorimetry
 - improving neutrino energy resolution

• Could provide timing to emulstion detector







Near emulsion detector

Advantages of the emulsion detector

- Can reconstruct all charged particle tracks with high precision
- Can detect gammas via conversion
- Good electron/muon/hadron discrimination

Disadvantages of the emulsion detectors

- No timing information
 - But can be restored by connecting tracks with FGD
- Price per mass
- No online event reconstruction
- Labour intensive



Water target emulsion detector

Photo from T. Fukuda of NINJA collaboration

Usage in ESSnuSB

- Study of neutrino interaction topology
- Measurement of interaction cross-section







Two lonely dots in our energy region

Figure 51.1: Measurements of per nucleon ν_{μ} and $\overline{\nu}_{\mu}$ CC inclusive scattering cross sections divided by neutrino energy as a function of neutrino energy. Note the transition between logarithmic and linear scales occurring at 100 GeV. Neutrino cross sections are typically twice as large as their corresponding antineutrino counterparts, although this difference can be larger at lower energies. NC cross sections (not shown) are generally smaller compared to the CC case.

Taken from PDG 2020: P.A. Zyla et al. (Particle Data Group), Prog. Theor. Exp. Phys. 2020, 083C01 (2020)

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Why measure cross-section?

• 4 different "tunes" of GENIE 4.0.6 produce significantly differen curves







Cross-section measurement

Main problem:

- Event rate (what we measure) is proportional to (flux) x (cross section)
- So, we need one to measure the other, if using event rate as observable

Strategies:

- Use elastic scattering of neutrinos on electrons (known cross section) to constrain the flux
 - measured in the Near WC detector
- Having constraint on the flux, we can measure interaction cross-sections in all Near Detectors:
 - WC, Super FGD, emulsion
- Use current **and future** data from other experiments (e.g. T2K)





Neutrino scattering on electrons

Basic idea:

- Cross section for neutrino electron scattering can be calculated to high precision (NLO seems to be enough for all practical purposes)
- Knowing a cross-section, one can measure the flux at Near Detector site
 - More precisely, put tight constraints on the flux

Main problem:

- Neutrino cross-section scales with target mass
 - having electron as a target, the cross-section is much smaller than having nucleon as a target

Event selection:

• v - e scattering has a very forward single electron in the final state







Near detectors conclusions strategy

- Main purposes:
 - Constrain the prompt neutrino flux
 - Measure neutrino interaction cross-sections (both inclusive and exclusive)
- Main difficulty:
 - To disentangle the cross-section from flux
- Strategy:
 - Constrain flux by measuring neutrino-electron scattering
 - Use the constrained flux to measure (inclusive and inclusive) crosssections for interactions of neutrinos with water





The End

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