

Analytical solutions for the free surface Euler & Navier-Stokes systems

Some ingredients for the validation of numerical schemes

M.-O. Bristeau, B. Di Martino, J. Sainte-Marie,
F. Souillé & S. Téchéné

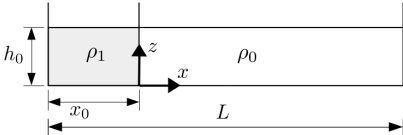
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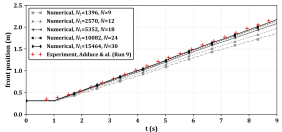
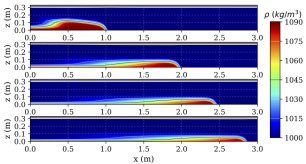
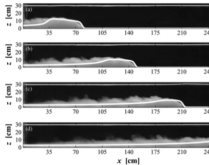
Commodore - January 2020 - Hambourg

A typical example

- The experimental device [Adduce *et al.*, J. Hydraul. Eng. 2012] (anim) (anim_low)

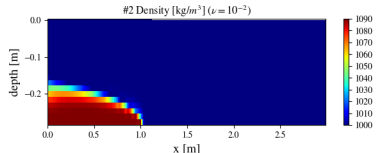
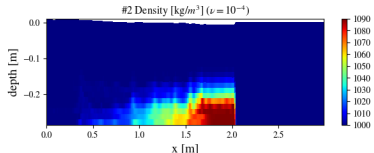
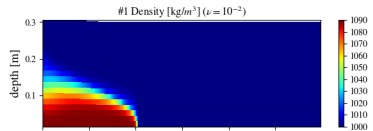
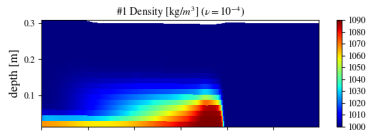


- Measurements versus simulation results

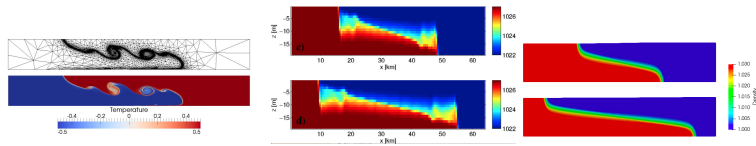


A typical example (cont'd)

Comparison of 2 different codes (one of the experiments proposed by Adduce)



Other codes



See also Wroniszewski et al. Benchmarking of Navier-Stokes codes for free surface simulations, Coastal Eng., 91:1-17,2014

Conclusion

- The experiment

- errors/uncertainties (initial conditions)
- non-hydrostatic / hydrostatic
- turbulence...

- The numerical codes

- important differences
- origins ?

- Analysis

- no existence result, out of reach
- shape of the interface ?

- Numerical analysis required

- consistency, domain invariant, order of the scheme, discrete entropy
- the two num. schemes are not consistent
- accuracy versus stability

- Use of “reference” solutions

- The model

$$\begin{cases} \dot{\rho} + \operatorname{div}(\rho \underline{\mathbf{u}}) = 0 \\ \frac{\dot{\rho}}{\rho} \underline{\mathbf{u}} + (\underline{\mathbf{u}} \cdot \nabla)(\rho \underline{\mathbf{u}}) + \nabla p = \rho \mathbf{G} \\ \frac{\dot{\rho}}{\rho} T + \operatorname{div}(\rho T \underline{\mathbf{u}}) = \mu \Delta T \\ \rho = \rho(T) \end{cases}$$

Contents

The analytical solutions

- Stationary / Non-stationary
- Euler / Navier-Stokes
- Hydrostatic / Non-hydrostatic
- Constant density / Variable density

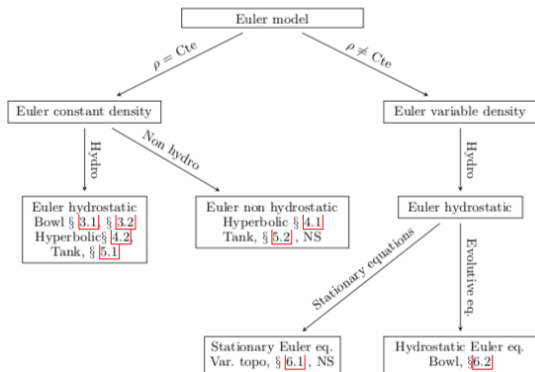
Allow to test

- Time discretization and/or space discretization
- Wet/dry interfaces
- Treatment boundary conditions
- Viscous terms discretization

Main idea

- Useful to have “reference” solutions when the numerical analysis is very complex
- Accuracy versus stability

The proposed analytical solutions



- Fully described in “Some analytical solutions for validation of free surface flow computational codes”
<https://hal.archives-ouvertes.fr/hal-01831622>
- New solutions or extensions of existing ones
- Also wave propagation

Shallow water analytical sol. of the Euler system

The 2d Euler system (u, w)

$$\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0$$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + w \frac{\partial u}{\partial z} + \frac{\partial p}{\partial x} = 0$$

$$\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + w \frac{\partial w}{\partial z} + \frac{\partial p}{\partial z} = -g$$

$$\frac{\partial \eta}{\partial t} + u_s \frac{\partial \eta}{\partial x} - w_s = 0, \quad u_b \frac{\partial z_b}{\partial x} - w_b = 0, \quad p|_s = p_0^a$$

The SW assumption $u(x, z, t) = \bar{u}(x, t)$

$$w = -(z - z_b) \frac{\partial \bar{u}}{\partial x} + \bar{u} \frac{\partial z_b}{\partial x}$$

Expression of the pressure

$$-z \left(\frac{\partial^2 \bar{u}}{\partial x \partial t} + \bar{u} \frac{\partial^2 \bar{u}}{\partial x^2} - \left(\frac{\partial \bar{u}}{\partial x} \right)^2 \right) + \frac{\partial^2 (z_b \bar{u})}{\partial x \partial t} + \bar{u} \frac{\partial^2 (z_b \bar{u})}{\partial x^2} - \frac{\partial \bar{u}}{\partial x} \frac{\partial (z_b \bar{u})}{\partial x} + \frac{\partial p}{\partial z} = -g$$

Hence

$$\frac{\partial^2 p}{\partial x \partial z} = \frac{\partial^3 p}{\partial z^3} = 0 \Rightarrow p = p_0^a + \frac{a(t)}{2} (\eta^2 - z^2) + b(t) (\eta - z)$$

Shallow water analytical sol. of the Euler system

The only SW solutions of the 3d Euler system

$$h(t, x, y) = \max \{0, \alpha f(t) - b_0 - z_b(x, y)\}$$

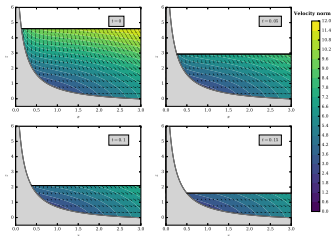
$$u(t, x, y, z) = f(t)(x \cos \theta + y \sin \theta + \beta) \cos \theta$$

$$v(t, x, y, z) = f(t)(x \cos \theta + y \sin \theta + \beta) \sin \theta$$

$$w(t, x, y, z) = -f(t)(z + b_0)$$

$$p(t, x, y, z) = p^{a,0}(t) + f^2(t)(\eta(x, y, t) - z^2) + (2b_0 f^2(t) + g)(\eta - z)$$

where $f(t) = 1/(t - t_0 + t_1)$, $z_b(x, y) = \frac{c_0}{x \cos \theta + y \sin \theta + \beta} - b_0$, and $p^{a,0}(t)$ a given function.



Animations

(With topography)

(Without topography)

also valid for Navier-Stokes

Shallow water sol. of the hydro. Euler system

The only SW sol. of the 3d hydrostatic Euler system

$$h(t, x, y) = \max \{0, \alpha f(t) - b_0 - zb(x, y)\}$$

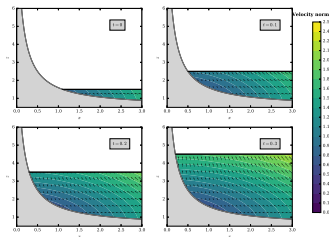
$$u(t, x, y, z) = -f(t)(x \cos \theta + y \sin \theta + \beta) \cos \theta$$

$$v(t, x, y, z) = -f(t)(x \cos \theta + y \sin \theta + \beta) \sin \theta$$

$$w(t, x, y, z) = f(t)(z + b_0)$$

$$p(t, x, y, z) = p^{a,0}(t) + p^{a,1}(x, t) + g(\eta - z)$$

where $f(t) = 1/(t - t_0 + t_1)$, $z_b(x, y) = \frac{c_0}{x \cos \theta + y \sin \theta + \beta} - b_0$, and $p^{a,0}(t)$ a given function.



The hydrostatic & incompressible 3d Euler system

- The governing equations (hydrostatic)

$$\begin{cases} \nabla \cdot \mathbf{U} = 0 \\ \rho_0 \left(\frac{\partial \mathbf{U}}{\partial t} + \nabla \cdot (\mathbf{U} \otimes \mathbf{U}) \right) + \nabla p = -\rho_0 \mathbf{g} \\ \frac{\partial \eta}{\partial t} + u_s \frac{\partial \eta}{\partial x} - w_s = 0, \quad u_b \frac{\partial z_b}{\partial x} - w_b = 0 \\ \rho|_s = \rho_0^a \end{cases}$$

- Analytical solution (anim)

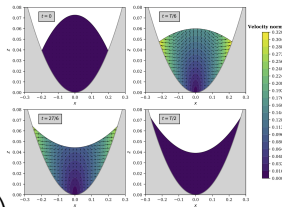
$$H(t, x, y) = \max \left\{ 0, \frac{1}{r^2} f \left(\frac{r^2}{\gamma \cos(\omega t) - 1} \right) \right\}$$

$$u(t, x, y, z) = x \left(\beta \left(z - z_b - \frac{H}{2} \right) + \frac{\omega \gamma \sin(\omega t)}{2(1 - \gamma \cos(\omega t))} \right)$$

$$v(t, x, y, z) = y \left(\beta \left(z - z_b - \frac{H}{2} \right) + \frac{\omega \gamma \sin(\omega t)}{2(1 - \gamma \cos(\omega t))} \right)$$

$$w(t, x, y, z) = -\frac{\partial}{\partial x} \int_{z_b}^z u dz - \frac{\partial}{\partial y} \int_{z_b}^z v dz$$

$$\rho(t, x, y, z) = \rho^{a,0} + g(H + z_b - z)$$



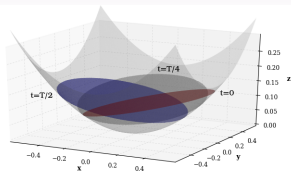
- Also moving bottom

(3d extension of Thacker)

The 3d Euler system with varying density

- Hydrostatic Euler system (Boussinesq)

$$\left\{ \begin{array}{l} \nabla \cdot \mathbf{U} = 0 \\ \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{U}) = 0 \\ \rho_0 \left(\frac{\partial \tilde{\mathbf{U}}}{\partial t} + \nabla \cdot (\tilde{\mathbf{U}} \otimes \mathbf{U}) \right) + \nabla p = \rho \mathbf{g} \\ \frac{\partial \eta}{\partial t} + u_s \frac{\partial \eta}{\partial x} - w_s = 0, \quad u_b \frac{\partial z_b}{\partial x} - w_b = 0 \\ \rho|_s = \rho_0^a \end{array} \right.$$



- Analytical solution (anim)

$$H(t, x, y) = \max \left\{ 0, h_0 - \alpha \frac{(x - \eta \cos(\omega t))^2 + (y - \eta \sin(\omega t))^2}{2} \right\}$$

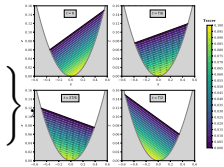
$$u(t, x, y, z) = -\eta \omega \sin(\omega t),$$

$$v(t, x, y, z) = \eta \omega \cos(\omega t),$$

$$w(t, x, y, z) = -\alpha \eta \omega (x \sin(\omega t) - y \cos(\omega t)),$$

$$\rho(t, x, y, z) = \rho^a(t) + \int_z^{H+z_b} \rho(\phi(t, x, y, z_1)) dz_1,$$

$$\phi(t, x, y, z) = a(H + z_b - z),$$



- For any nonnegative function $\rho : s \mapsto \rho(s)$

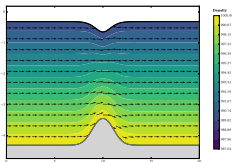
Remarks:

- ◇ ρ can be discontinuous (slow convergence of the num. scheme)
- ◇ Numerical stability of contact discontinuities?

The 3d Euler system with varying density

- Stationary hydrostatic Euler system (Boussinesq)

$$\left\{ \begin{array}{l} \nabla \cdot \mathbf{U} = 0 \\ \nabla \cdot (\rho \mathbf{U}) = 0 \\ \rho_0 \nabla \cdot (\tilde{\mathbf{U}} \otimes \mathbf{U}) + \nabla p = \rho \mathbf{g} \\ u_s \frac{\partial \eta}{\partial x} - w_s = 0, \quad u_b \frac{\partial z_b}{\partial x} - w_b = 0 \\ \rho|_s = \rho_0^a \end{array} \right.$$

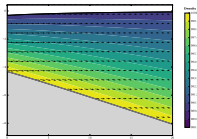


- Analytical solution

$$u(x, z) = f_1(x) + \frac{g\alpha}{2\rho_0} z^2$$

$$w(x, z) = -f_1'(x)(z - z_b(x)) + f_1(x)z_b'(x) + \frac{\alpha g}{2\rho_0} z_b^2(x)z_b'(x)$$

$$\rho(x, z) = \rho_0 - \alpha f_1(x)(z - z_b(x)) - \frac{\alpha^2 g}{6\rho_0} (z^3 - z_b^3(x))$$



with the function $f_1(x)$ defined by

$$f_1(x) = \frac{\alpha g}{2\rho_0} (h^2(x) - z_b^2(x)) - \frac{1}{\sqrt{3}\rho_0} \sqrt{(\alpha^2 g^2 h^3(x) - 6\rho_0^2 g)(h(x) + z_b(x)) - 6C_0}$$

A special function

- The **LambertW** function

$$\mathcal{LambertW}(x)e^{\mathcal{LambertW}(x)} = x$$

- In other words $\mathcal{LambertW}(\cos(x))$ is the function $y = y(x)$ defined by

$$ye^y = \cos(x)$$

- The equation

$$ye^y = a \cos(x)$$

admits several complex solutions, a **unique** real solution when $a \leq \frac{1}{e}$

- For small a

$$\mathcal{LambertW}(a \cos(x)) = a \cos(x) + \mathcal{O}(a^2)$$

Analytical solutions: linearized Euler

Proposition (with an energy balance)

Let f the function defined by

$$f(x) = -\frac{n}{k} \text{LambertW}\left(a \cos\left(\frac{x}{n} + b\right)\right)$$

where $|a| < 1/e$, $b, k > 0$ and $n > 0$ are constants.

Up to terms in $\mathcal{O}(e^{-\frac{H_0 k}{n}})$, the functions H , u , w and p defined by

$$H = H_0 + f(kx - \omega t)$$

$$u = \frac{\omega}{n} e^{\frac{k}{n}(z-H)} f(kx - \omega t)$$

$$w = -\frac{n}{k} \frac{\partial u}{\partial x}$$

$$p = p^a(t) + g(H_0 - z) - \frac{gan}{k} e^{\frac{k}{n}(z-H_0)} \cos\left(\frac{kx - \omega t}{n} + b\right)$$

are analytical solutions of the linearized Euler system iff $\frac{\omega}{k} = \sqrt{\frac{gn}{k}}$.

Analytical solutions: Euler

Proposition

Let f the function defined by

$$f(x) = -\frac{n}{k} \text{LambertW}\left(a \cos\left(\frac{x}{n} + b\right)\right)$$

where $|a| < 1/e$, $b, k > 0$ and $n > 0$ are constants.

Up to terms in $\mathcal{O}(e^{-\frac{H_0 k}{n}})$, the functions H , u , w and p defined by

$$H = H_0 + f(kx - \omega t), \quad u = \frac{\omega}{n} e^{\frac{k}{n}(z-H)} f(kx - \omega t), \quad w = -\frac{n}{k} \frac{\partial u}{\partial x}$$

$$p = p^a(x, t) + g(H_0 - z) - \frac{ga^2 n}{2k} e^{\frac{2k}{n}(z-H_0)} \\ - \frac{gan}{k} e^{\frac{k}{n}(z-H_0)} \cos\left(\frac{kx - \omega t}{n} + b\right)$$

$$p^a(x, t) = -\frac{ga^2 n}{2k} e^{\frac{2k}{n}(H-H_0)}$$

are analytical solutions of the Euler system iff $\frac{\omega}{k} = \sqrt{\frac{gn}{k}}$.

Standing waves

- Basic idea

$$\cos\left(\frac{kx - \omega t}{n}\right) + \cos\left(\frac{kx + \omega t}{n}\right) = 2 \cos\left(\frac{2kx}{n}\right) \cos\left(\frac{2\omega t}{n}\right)$$

- (anim)

Proposition

With $f(x) = -\frac{n}{k} \mathcal{L}ambertW\left(a \cos\left(\frac{x}{n}\right)\right)$, up to terms in

$\mathcal{O}\left(e^{-\frac{H_0 k}{n}}, a^3\right)$, the functions

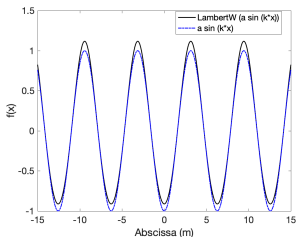
$$H = H_0 + f(kx - \omega t) + f(kx + \omega t) + 2 \frac{a^2 n}{k} \cos^2\left(\frac{kx}{n}\right)$$

$$u = \frac{\omega}{n} e^{\frac{k}{n}(z-H)} \left(f(kx - \omega t) + f(kx + \omega t) + 2 \frac{a^2 n}{k} \cos^2\left(\frac{kx}{n}\right) \right)$$

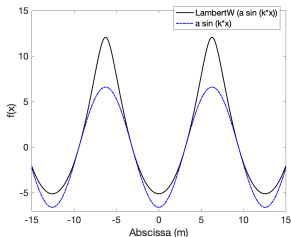
...

are analytical solutions of the Euler system, linearized Euler system

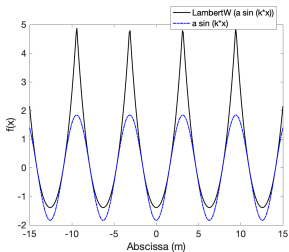
$$f(x) = -\frac{n}{k} \text{LambertW}\left(a \cos\left(\frac{x}{n} + b\right)\right)$$



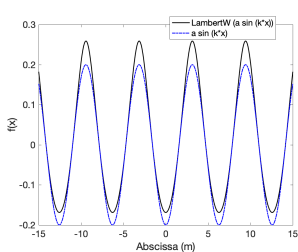
$$a = 0.1 \text{ m}, k = 0.1 \text{ m}^{-1}, n = 1$$



$$a = 0.33 \text{ m}, k = 0.1 \text{ m}^{-1}, n = 2$$



$$a = e^{-1} \text{ m}, k = 0.2 \text{ m}^{-1}, n = 1$$



$$a = 0.2 \text{ m}, k = 1 \text{ m}^{-1}, n = 1$$

Conclusion: about these analytical solutions

- Accessible by simulations prescribing
 - geometrical domain
 - initial and boundary conditions
- Stability
 - some of them also valid for Navier-Stokes (viscous fluid)
 - but not enough, analysis result required
 - sometimes convergence very low (variable density case)
- Should be completed with “reference solutions”